## On reversible solutions of SPDEs PIMS 4/7/2005

Aim: study scaling limits of evolving interfaces (see [Giacomin,Lebowitz,Presutti] or Funaki's St-Flour Lectures)

Equilibrium: Gibbs measure, RW model for interface:  $V \in C^2(\mathbb{R})$  even

$$S_n = X_1 + \dots + X_n, \qquad (X_i)_i \quad \text{IID}$$

$$X_i \sim \frac{1}{Z} \exp(-V(r)) \, dr, \quad 0 < c_- < V'' < c_+ < \infty$$

$$\mathbb{E}[X_i] = 0, \qquad \mathbb{E}[X_i^2] = 1$$
Then the interface is  $\{(n, S_n) : n \leq N + 1\}$ 
conditioned on

$$\{S_1, \ldots, S_N \ge 0, S_0 = S_{N+1} = 0\}$$

We call  $\mu_N^+$  the law:

$$\mu_N^+(d\phi) = \frac{1}{Z_N^+} \mathbf{1}_{(\phi \ge 0)} e^{-\sum_i V(\phi_{i+1} - \phi_i)} d\phi$$
  
ith  $\phi_0 = \phi_{N+1} = 0.$ 

Under Brownian rescaling:

$$Y_t = \frac{1}{\sqrt{N}} S_{\lfloor Nt \rfloor} \implies e_t$$

where e is the normalized Brownian excursion. Notation:

$$\nu_N := Y \circ \mu_N^+ \implies \nu$$

law of  $\boldsymbol{e}$ 

W

Natural reversible dynamics: 
$$\phi_t \in \mathbb{R}^N_+$$
,  
 $(\partial \phi)_i := \phi_{i+1} - \phi_i$ ,  $(\partial^* \phi)_i := \phi_i - \phi_{i-1}$   
 $\begin{cases} d\phi_i = \frac{1}{2} \{ \partial V'(\partial^* \phi) \}_i dt + dw_i + dl_i \} \\ \phi_0(t) = \phi_{N+1}(t) = 0, \\ \phi_i \ge 0, \quad dl_i \ge 0, \quad \int_0^\infty \phi_i(t) \ dl_i(t) = 0 \end{cases}$ 

 $\exists$ ! stationary  $\phi$ . Funaki-Olla [SPA 01]:

$$\Phi_N(t,x) := \frac{1}{\sqrt{N}} \sum_i \phi_x(N^2 t) \, \mathbf{1}_{\left[\frac{i-1}{N}, \frac{i}{N}\right]}(x)$$

 $\Longrightarrow \Phi_N \rightarrow$  unique stationary solution of ...

Nualart and Pardoux [PTRF 92]:  $\exists !(u, \eta)$ :

$$\begin{cases} \frac{\partial u}{\partial t} = \frac{1}{2} \frac{\partial^2 u}{\partial x^2} + \dot{W} + \eta(t, x) \\ u(0, x) = u_0(x), \quad x \in [0, 1] \\ u(t, 0) = u(t, 1) = 0, \quad t \ge 0 \\ u \ge 0, \ d\eta \ge 0, \ \int u \, d\eta = 0 \end{cases}$$

- $\dot{W}$  space-time white noise
- $\eta$  is a reflecting measure on  $\mathbb{R}_+ \times [0, 1]$

Existence: by penalization

$$\varepsilon > 0$$
:  $\frac{\partial u_{\varepsilon}}{\partial t} = \frac{1}{2} \frac{\partial^2 u_{\varepsilon}}{\partial x^2} + \dot{W} + \frac{(u_{\varepsilon})^-}{\varepsilon}$ 

Then  $u_{\varepsilon} \uparrow u$  uniformly as  $\varepsilon \downarrow 0$ .

Results on the contact set  $\{(t, x) : u(t, x)\}$ .

New proof of Funaki-Olla's result, based on three main properties:

- 1. trivial convergence of Dirichlet forms
- 2. uniform strong Feller property for  $\Phi_N$ :

$$|P_t^N F(\Phi) - P_t^N F(\Phi')| \leq \frac{\|F\|_{\infty}}{t^{1/2}} \|\Phi - \Phi'\|_{L^2(0,1)}$$
$$P_t^N F(\Phi) := \mathbb{E}_{\Phi}[F(\Phi_N(t))]$$

3. convergence of integration by parts formulae (IbPF) for  $\nu_N$  to IbPF for  $\nu$  Reflecting Brownian Motion:

$$dX = dB + dL,$$
  
$$X \ge 0, \quad dL \ge 0, \quad \int_0^\infty X \, dL = 0$$

associated to the Dirichlet form:

$$D(\varphi,\psi) = \frac{1}{2} \int_0^\infty \varphi' \,\psi' \,dx.$$

Informally: 
$$L_t = \frac{1}{2} \int_0^t \delta_0(X_s) \, ds$$
. IbPF:  
 $\int_0^\infty \varphi' \, dx = -\varphi(0) = -\delta_0(\varphi), \qquad \varphi \in C_c^1(\mathbb{R})$ 

The Dirac Delta is the Revuz measure of L.

Back to interfaces: now we want the dynamics to be conservative, i.e.

$$\sum_{i} \phi_i(t) = \sum_{i} \phi_i(0) \qquad \forall t \ge 0$$

(constant droplet volume). The natural dynamics is:

$$\begin{cases} d\phi = -\frac{1}{2}\partial\partial^* \{\partial V'(\partial^*\phi) \, dt + dl_i\} + \partial dw \\ \phi_0(t) = \phi_{N+1}(t) = 0, \\ \phi_i \ge 0, \quad dl_i \ge 0, \quad \int_0^\infty \phi_i(t) \, dl_i(t) = 0 \end{cases}$$

 $\exists !$  stationary  $\phi$  after fixing the droplet volume.

$$\Phi_N(t,x) := \frac{1}{\sqrt{N}} \sum_i \phi_x(N^4 t) \,\mathbf{1}_{\left[\frac{i-1}{N},\frac{i}{N}\right)}(x)$$

 $\Phi_N \rightarrow$  unique stationary solution of  $\ldots$ 

Stochastic Cahn-Hilliard equation: (joint with A.Debussche)

$$\begin{cases} \frac{\partial u}{\partial t} = -\frac{1}{2} \frac{\partial^2}{\partial x^2} \left( \frac{\partial^2 u}{\partial x^2} + \eta(t, x) \right) + \frac{\partial}{\partial x} \dot{W} \\ u(t, 0) = u(t, 1) = 0, \\ \partial^3 u(t, 0) = \partial^3 u(t, 1) = 0, \\ u \ge 0, \ d\eta \ge 0, \ \int u \, d\eta = 0 \end{cases}$$

Difficulty even for existence: try penalization

$$\frac{\partial u^{\epsilon}}{\partial t} = -\frac{1}{2} \frac{\partial^2}{\partial x^2} \left( \frac{\partial^2 u^{\epsilon}}{\partial x^2} + \frac{(u^{\epsilon})^-}{\epsilon} \right) + \frac{\partial}{\partial x} \dot{W}$$

no monotonicity; tightness trivial known only for stationary solutions

the 1.-2.-3. proof above gives convergence of the semigroups and Strong Feller for the limit, and this is enough to conclude Also, the 1.-2.-3. proof gives convergence of the stationary interface.

at present, we have a problem with uniqueness for CH: we expect pathwise uniqueness (easy with different boundary conditions, difficult here because the mass of  $\eta$  is expected to be  $+\infty$ )

the Cahn-Hilliard equation is a gradient system in  $H^{-1}(0,1)$ : in this norm localization works very badly

Last model: at equilibrium we have pinning, i.e. the interface gets a reward  $\epsilon > 0$  every time it touches the wall. the measure is  $\mu_{\epsilon,N}^+$ , the Brownian rescaling  $\nu_{\epsilon,N}$ 

by tuning  $\epsilon$ , we have a phase transition (joint with Deuschel, Giacomin):

•  $\epsilon > \epsilon_c$ : convergence to flat interface  $\equiv 0$ 

• 
$$\epsilon < \epsilon_c$$
:  $\nu_{\epsilon,N} \Longrightarrow \nu$  (as if  $\epsilon = 0$ )

•  $\epsilon = \epsilon_c$ :  $\nu_{\epsilon,N} \Longrightarrow$  law of reflecting Brownian bridge

The natural dynamics has **sticky** reflection:

$$\begin{cases} d\phi_i = \mathbf{1}_{(\phi_i(t)>0)} \left[ \frac{1}{2} \{ \partial V'(\partial^* \phi) \}_i dt + dw_i \right] \\ + \frac{1}{2\epsilon} \mathbf{1}_{(\phi_i(t)=0)} dt \\ \phi_0(t) = \phi_{N+1}(t) = 0, \end{cases}$$

 $\exists$  stationary  $\phi$  (unique?).

$$\Phi_N(t,x) := \frac{1}{\sqrt{N}} \sum_i \phi_x(N^2 t) \, \mathbf{1}_{\left[\frac{i-1}{N}, \frac{i}{N}\right]}(x)$$

at  $\epsilon = \epsilon_c$ ,  $\Longrightarrow \Phi_N \rightarrow$  unique stationary solution of ...

$$\begin{cases} \frac{\partial u}{\partial t} = \frac{1}{2} \frac{\partial^2 u}{\partial x^2} + \dot{W} - \left| \frac{\partial u}{\partial x} \right|^2 : \eta(t, x) \\ u(0, x) = u_0(x), \quad x \in [0, 1] \\ u(t, 0) = u(t, 1) = 0, \quad t \ge 0 \\ u \ge 0, \ d\eta \ge 0, \ \int u \, d\eta = 0 \\ : \left| \frac{\partial u}{\partial x} \right|^2 := \lim_{\epsilon} \left[ \left| \frac{\partial u_{\epsilon}}{\partial x} \right|^2 - C_{\epsilon} \right] \end{cases}$$

This renormalized term reminds of the KPZ equation, but this one has **reversible** solutions!

Done: IbPF for the reflecting Brownian motion